

On the second mixed moment of the characteristic polynomials of sparse hermitian random matrices

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Outline

- 1 Basic concepts of the random matrix theory
- 2 The mixed moments of the characteristic polynomials
- 3 The ensemble of the sparse random matrices
- 4 The Grassmann variables method
- 5 Formulations of the theorems

Random matrices

Definition

The sequence M_n of the $n \times n$ matrices, which entries are random variables, is called a random matrix ensemble.

Examples

- Wigner ensemble

$$M^T = M \text{ or } M^* = M$$

$$M_{jk} = n^{-1/2} w_{jk},$$

$$\{w_{jk}\} - \text{i.i.d.}, \quad E\{w_{jk}\} = 0, \quad E\{|w_{jk}|^2\} = 1$$

- The adjacency matrices of random graphs

$$M_{jk} = \begin{cases} 1 & \text{with probability } \frac{p_n}{n}; \\ 0 & \text{with probability } 1 - \frac{p_n}{n}. \end{cases}$$

Global regime

Normalized counting measure of eigenvalues (NCM) and linear eigenvalue statistics

$$N_n(\Delta) = \frac{1}{n} \sum_{j=1}^n \mathbb{1}_{\Delta}(\lambda_j^{(n)}), \quad \mathcal{N}_n[\varphi] = \sum_{j=1}^n \varphi(\lambda_j^{(n)}) = \text{tr } \varphi(M)$$

Questions

- 1 $N_n(\Delta) \xrightarrow{w} N(\Delta)$
- 2 Central Limit Theorem for linear eigenvalue statistics

$\sigma = \text{supp } N$ is called the spectrum.

For the Wigner ensemble

$$N(\Delta) = \int_{\Delta} \rho_{\text{sc}}(\lambda) d\lambda, \quad \rho_{\text{sc}}(\lambda) = \frac{1}{2\pi} \sqrt{4 - \lambda^2} \cdot \mathbb{1}_{[-2,2]}, \quad \sigma = [-2, 2].$$

Local regime

The spectral correlation functions are

$$p_k^{(n)}(\lambda_1, \dots, \lambda_k) = \int p_n(\lambda_1, \dots, \lambda_n) d\lambda_{k+1} \dots d\lambda_n.$$

where $p_n(\lambda_1, \dots, \lambda_n)$ is the joint probability density of the eigenvalues.

Dyson universality conjecture

$$\begin{aligned} \lim_{n \rightarrow \infty} (\rho_n(\lambda_0))^{-k} p_k^{(n)}(\lambda_0 + x_1/n\rho_n(\lambda_0), \dots, \lambda_0 + x_k/n\rho_n(\lambda_0)) \\ = \det \left\{ \frac{\sin \pi(x_i - x_j)}{\pi(x_i - x_j)} \right\}_{i,j=1}^k \end{aligned}$$

where $\rho_n(\lambda) = p_1^{(n)}(\lambda)$.

The mixed moments of the characteristic polynomials

Let M_n be some ensemble of random matrices. Consider the second mixed moment or the correlation function of characteristic polynomials

$$F_2(\Lambda) = E \{ \det(M_n - \lambda_1) \det(M_n - \lambda_2) \}, \quad \Lambda = \text{diag}\{\lambda_j\}_{j=1}^2$$

$$\text{a) } \lambda_j = \lambda_0 + \frac{x_j}{n}, \quad \text{b) } \lambda_j = \lambda_0 + \frac{x_j}{n^{2/3}}$$

$$F_2(\Lambda) \xrightarrow[n \rightarrow \infty]{} ?$$

Some results

- Keating, Snaith (2000)
- Brezin, Hikami (2000, 2001)
- Strahov, Fyodorov (2002, 2003); Fyodorov, Khoruzhenko (2006)
- Götze, Kösters (2008, 2009)
- T. Shcherbina (2011, 2013, 2014, 2015)

Sparse random matrices

Ensemble of the sparse hermitian random matrices

$$M_n = (d_{jk} w_{jk})_{j,k=1}^n,$$

where

$$d_{jk} = p^{-1/2} \begin{cases} 1 & \text{with probability } \frac{p}{n}; \\ 0 & \text{with probability } 1 - \frac{p}{n}. \end{cases}$$

and $\Re w_{jk}$, $\Im w_{jk}$, w_{ll} are independent Gaussian random variables with zero mean such that

$$E\{|w_{jk}|^2\} = 1.$$

The ensemble is studied in two regimes

- 1 $\lim_{n \rightarrow \infty} p < \infty$
- 2 $\lim_{n \rightarrow \infty} p = \infty$

The results on the ensemble of sparse random matrices

Global regime

- For $p \rightarrow \infty$ the normalized counting measure is the same as for the Wigner ensemble.
 - ▶ Rodgers, Bray (1988) on physical level of rigour;
 - ▶ Khorunzhy, Khoruzhenko, Pastur and M. Shcherbina (1992).
- For finite p the convergence of the normalized counting measure was proven by
 - ▶ Rodgers, De Dominicis (1990) on physical level of rigour;
 - ▶ Bauer, Golinelli (2001) for $w_{jk} = 1$;
 - ▶ Khorunzhy, M. Shcherbina, Vengerovsky (2004) in general case.
- Central Limit Theorem for linear eigenvalue statistics was proven by M. Shcherbina, Tirozzi for finite p (2010) and for $p \rightarrow \infty$ (2012).

The results on the ensemble of sparse random matrices

Local regime

- In the papers by Erdős, Knowles, Yau, Yin (2012) and Huang, Landon, Yau (2015) it was rigorously proved that for $p \gg n^\varepsilon$ the spectral correlation functions converge in weak sense to that for Wigner ensemble.
- The conjecture of existing of the critical value $p_c > 1$ at which the correlation of eigenvalues is changed.
 - ▶ Evangelou, Economou (1992);
 - ▶ Fyodorov, Mirlin (1991, on the physical level of rigour).

Grassmann variables

Let $\{\psi_j, \bar{\psi}_j\}_{j=1}^n$ be a set of anticommuting variables, i.e.

$$\psi_j \psi_k + \psi_k \psi_j = \bar{\psi}_j \psi_k + \psi_k \bar{\psi}_j = \bar{\psi}_j \bar{\psi}_k + \bar{\psi}_k \bar{\psi}_j = 0.$$

In particular, $\psi_j^2 = \bar{\psi}_k^2 = 0$. The set generates a graded algebra \mathcal{A} of polynomials of $\{\psi_j, \bar{\psi}_j\}$, which is called the Grassmann algebra.

For an analytical function f it's domain can be extended to Grassmann algebra by following.

$$f(\chi + z_0) = \sum_{j=0}^{\infty} \frac{f^{(j)}(z_0)}{j!} \chi^j,$$

where χ is a polynomial of $\{\psi_j, \bar{\psi}_j\}$ with zero free term.

Integration over the Grassmann variables

The integral over the Grassmann variables is a linear functional, defined on the basis by the relations

$$\int d\psi_j = \int d\bar{\psi}_k = 0, \quad \int \psi_j d\psi_j = \int \bar{\psi}_k d\bar{\psi}_k = 1.$$

A multiple integral is defined to be the repeated integral. Moreover “differentials” $\{d\psi_j, d\bar{\psi}_j\}_{j=1}^n$ anticommute with each other and with $\{\psi_j, \bar{\psi}_j\}_{j=1}^n$.

For example, for a function f

$$f(\psi_1, \dots, \psi_n) = a_0 + \sum_{j=1}^n a_j \psi_j + \dots + a_{1, \dots, n} \prod_{j=1}^n \psi_j$$

we have by definition

$$\int f(\psi_1, \dots, \psi_n) d\psi_n \dots d\psi_1 = a_{1, \dots, n}.$$

Integration over the Grassmann variables

Let A be a positive definite $n \times n$ matrix. The following Gaussian integral is well-known

$$\frac{1}{\pi^n} \int \exp \left\{ - \sum_{j,k=1}^n \bar{z}_j A_{jk} z_k \right\} \prod_{j=1}^n d\Re z_j d\Im z_j = \frac{1}{\det A}.$$

The important analogue of this formula in Grassmann variables theory

$$\int \exp \left\{ - \sum_{j,k=1}^n \bar{\psi}_j A_{jk} \psi_k \right\} \prod_{j=1}^n d\bar{\psi}_j d\psi_j = \det A \quad (1)$$

is valid for any matrix A .

If A is a hermitian matrix with i.i.d. Gaussian entries then the l.h.s. of (1) can be easily averaged

$$E\{\det A\} = \int \exp \left\{ \sum_{j < k} \bar{\psi}_j \psi_k \bar{\psi}_k \psi_j \right\} \prod_{j=1}^n d\bar{\psi}_j d\psi_j$$

The derivation of the integral representation

$$\begin{aligned} F_2(\Lambda) &= E \{ \det(M_n - \lambda_1) \det(M_n - \lambda_2) \} \\ &= C \int \exp \{ \Phi ((\bar{\psi}_1, \psi_1), (\bar{\psi}_1, \psi_2), (\bar{\psi}_2, \psi_1), (\bar{\psi}_2, \psi_2)) \} d\Psi, \end{aligned}$$

where Φ is an even polynomial of the 4th degree. Using the Hubbard-Stratonovich transformation

$$\begin{aligned} e^{y^2} &= \frac{a}{\sqrt{\pi}} \int e^{2axy - a^2 x^2} dx, \\ e^{y^t} &= \frac{a^2}{\pi} \int e^{ay(u+iv) + at(u-iv) - a^2 u^2 - a^2 v^2} dudv \end{aligned}$$

we return to the usual integral representation

$$F_2(\Lambda) = C \iint \prod_j \exp \left\{ -\frac{1}{2} \text{tr} Q^2 + g(\bar{\psi}_{j1}, \psi_{j1}, \bar{\psi}_{j2}, \psi_{j2}, Q) \right\} dQ d\Psi = C \int \dots dQ,$$

where Q is 2×2 hermitian matrix.

Final integral representation

$$F_2(\Lambda) = C_n(X) \frac{ie^{\lambda_0(x_1+x_2)}}{x_1 - x_2} \int_{\mathbb{R}^3} (t_1 - t_2) \exp \left\{ -i \sum_{j=1}^2 x_j t_j \right\} e^{nf(t_1, t_2, s)} dt_1 dt_2 ds,$$

where

$$f(t_1, t_2, s) = \log \left(s \sqrt{\frac{2(n-p)}{np}} - t_1 t_2 \right) - \frac{1}{2} \left(\sum_{j=1}^2 (t_j + i\lambda_0)^2 + s^2 \right).$$

The second order correlation function

Theorem 1 [A.:16 (published in JSP)]

Consider the normalized second order correlation function

$$D_2(\Lambda) = \frac{F_2(\Lambda)}{\sqrt{F_2(\lambda_1 I)F_2(\lambda_2 I)}}.$$

Then we have for finite p

(i) for $p > 2$

$$\lim_{n \rightarrow \infty} D_2(\Lambda) = \begin{cases} \frac{\sin((x_1 - x_2)\sqrt{(\lambda_*^2 - \lambda_0^2)/2})}{(x_1 - x_2)\sqrt{(\lambda_*^2 - \lambda_0^2)/2}}, & \text{if } |\lambda_0| < \lambda_*, \\ 1, & \text{if } |\lambda_0| \geq \lambda_*; \end{cases}$$

$$\text{with } \lambda_* = \sqrt{\left(4 - \frac{8}{p}\right)_+}.$$

(ii) for $p \leq 2$

$$\lim_{n \rightarrow \infty} D_2(\Lambda) = 1,$$

where $\Lambda = \text{diag}\{\lambda_1, \lambda_2\} = \text{diag}\left\{\lambda_0 + \frac{x_1}{n}, \lambda_0 + \frac{x_2}{n}\right\}$, $\lambda_0, x_1, x_2 \in \mathbb{R}$.

The second order correlation function at the edge of the spectrum

Theorem 2 [A.:16 (published in JSP)]

Let $p \rightarrow \infty$ and $\lambda_0 = 2$. Then

(i) for $\frac{n^{2/3}}{p} \rightarrow \infty$

$$\lim_{n \rightarrow \infty} D_2 \left(2I + X/n^{2/3} \right) = 1;$$

(ii) for $\frac{n^{2/3}}{p} \rightarrow c$

$$\lim_{n \rightarrow \infty} D_2 \left(2I + X/n^{2/3} \right) = \frac{\mathbb{A}(x_1 + 2c, x_2 + 2c)}{\sqrt{\mathbb{A}(x_1 + 2c, x_1 + 2c)\mathbb{A}(x_2 + 2c, x_2 + 2c)}},$$

where $X = \text{diag}\{x_1, x_2\}$, $\mathbb{A}(x, y) = \frac{\text{Ai}(x)\text{Ai}'(y) - \text{Ai}'(x)\text{Ai}(y)}{x-y}$ and $\text{Ai}(x)$ is Airy function.

The correlation functions of higher order

Theorem 3 [A.:16 (published in JSP)]

Let $p \rightarrow \infty$, $\lambda_0 \in (-2, 2)$. Then

$$\lim_{n \rightarrow \infty} \frac{F_{2m}(\Lambda)}{\left(\prod_{j=1}^{2m} F_{2m}(\lambda_j I) \right)^{\frac{1}{2m}}} = \frac{\hat{S}_{2m}(X)}{\hat{S}_{2m}(I)},$$

where

$$\hat{S}_{2m}(X) = \frac{\det \left\{ \frac{\sin(\pi \rho_{sc}(\lambda_0)(x_j - x_{m+k}))}{\pi \rho_{sc}(\lambda_0)(x_j - x_{m+k})} \right\}_{j,k=1}^m}{\Delta(x_1, \dots, x_m) \Delta(x_{m+1}, \dots, x_{2m})}.$$